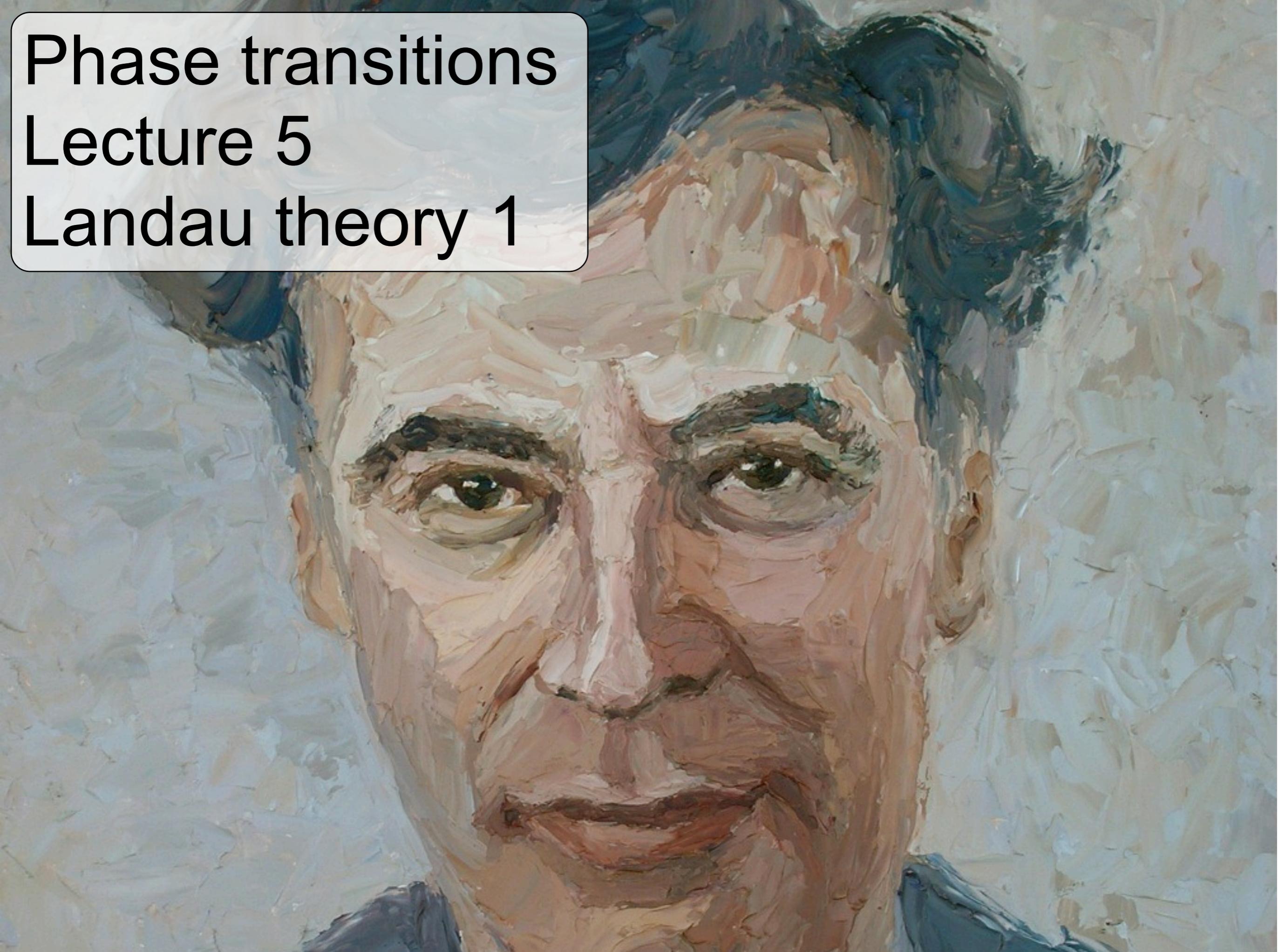


Phase transitions  
Lecture 5  
Landau theory 1



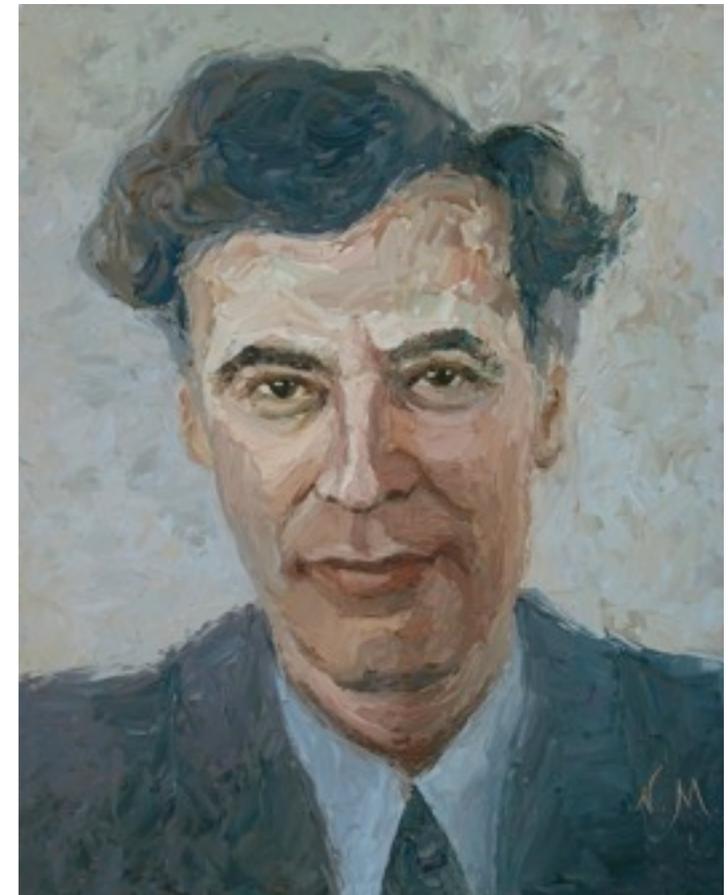


# Learning outcomes

- ▶ Know how the order parameter can be incorporated into a thermodynamic theory of phase transitions
- ▶ Understand the general form of the free energy function and how it changes through the phase transition
- ▶ Derive equations for the temperature-dependence of the order parameter

# Lev Landau

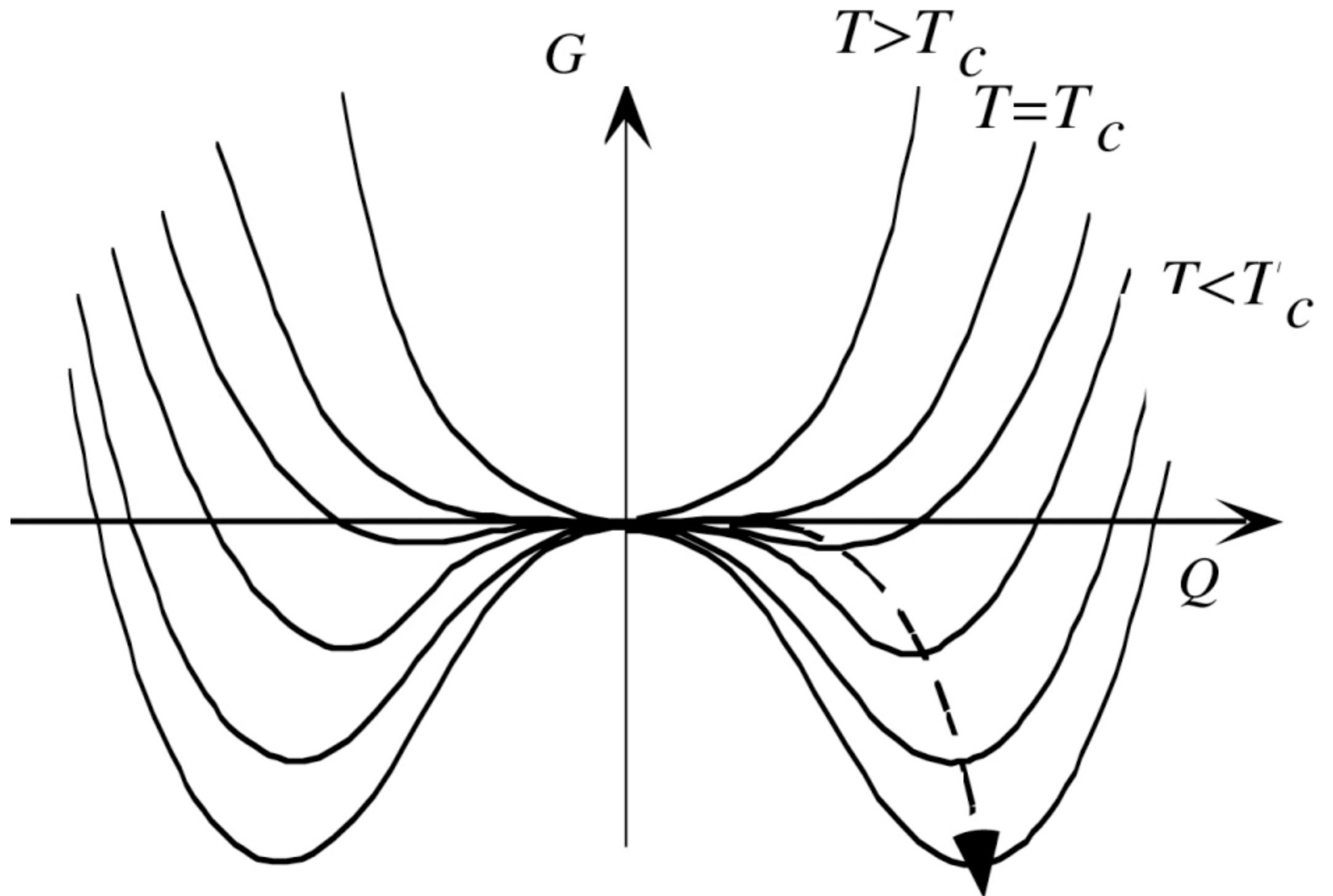
- ▶ 1908–1968
- ▶ Russian theoretician
- ▶ Nobel prize 1962
- ▶ First to identify the role of the order parameter and to develop a general theory of phase transitions



# Role of free energy

- ▶ The minimum of the free energy defines the equilibrium state
- ▶ Free energy varies with temperature; assume that the minimum point also changes with temperature
- ▶ Assume we can write the free energy as a function of order parameter

# Free energy for a second-order phase transition

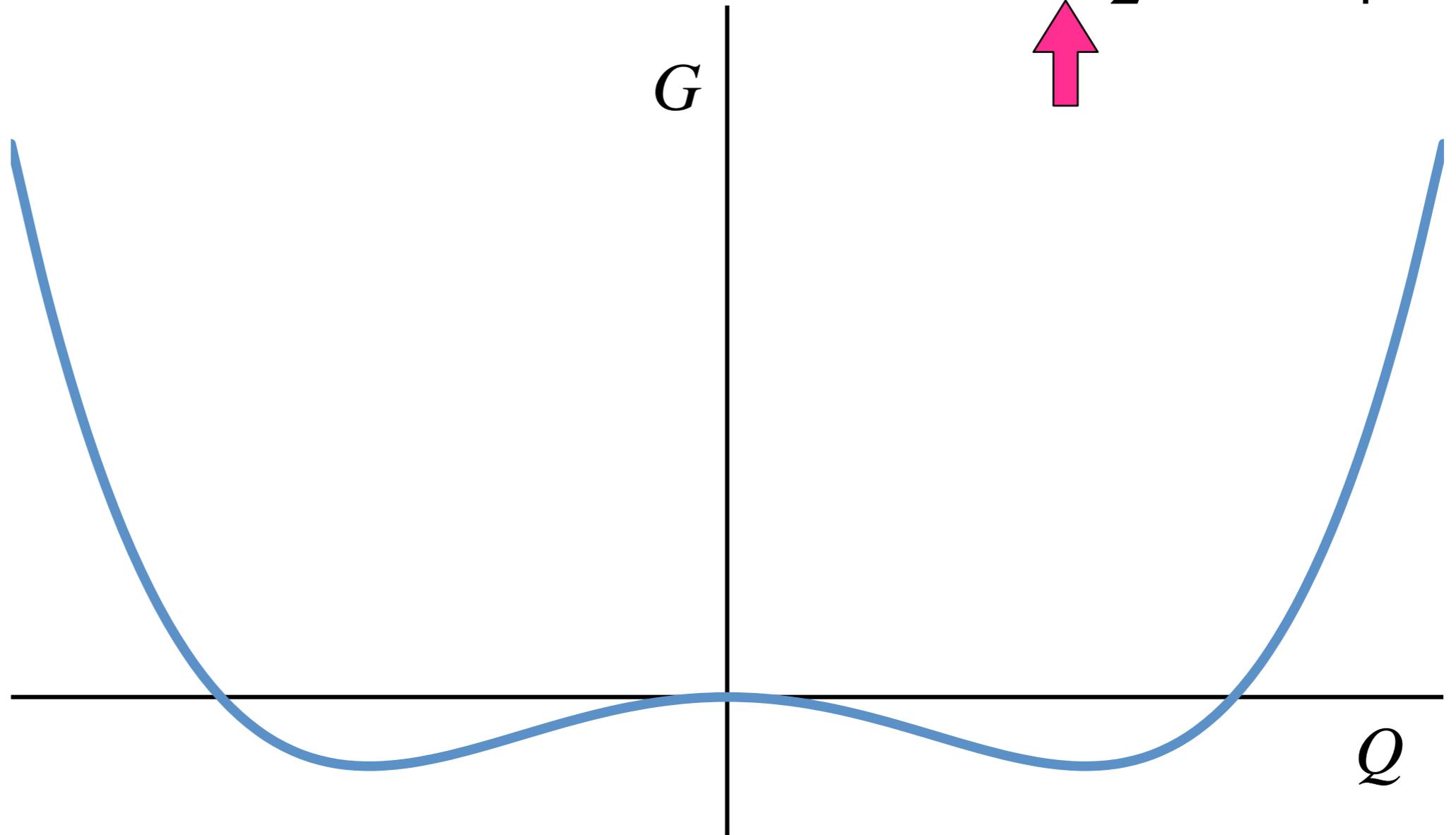


# Form of the free energy

- ▶ Assume we can expand  $G$  as a power series in  $Q$
- ▶ The condition  $G(-Q) = G(Q)$  requires that the power series only has even powers of  $Q$
- ▶ We will truncate the power series as soon as practically possible

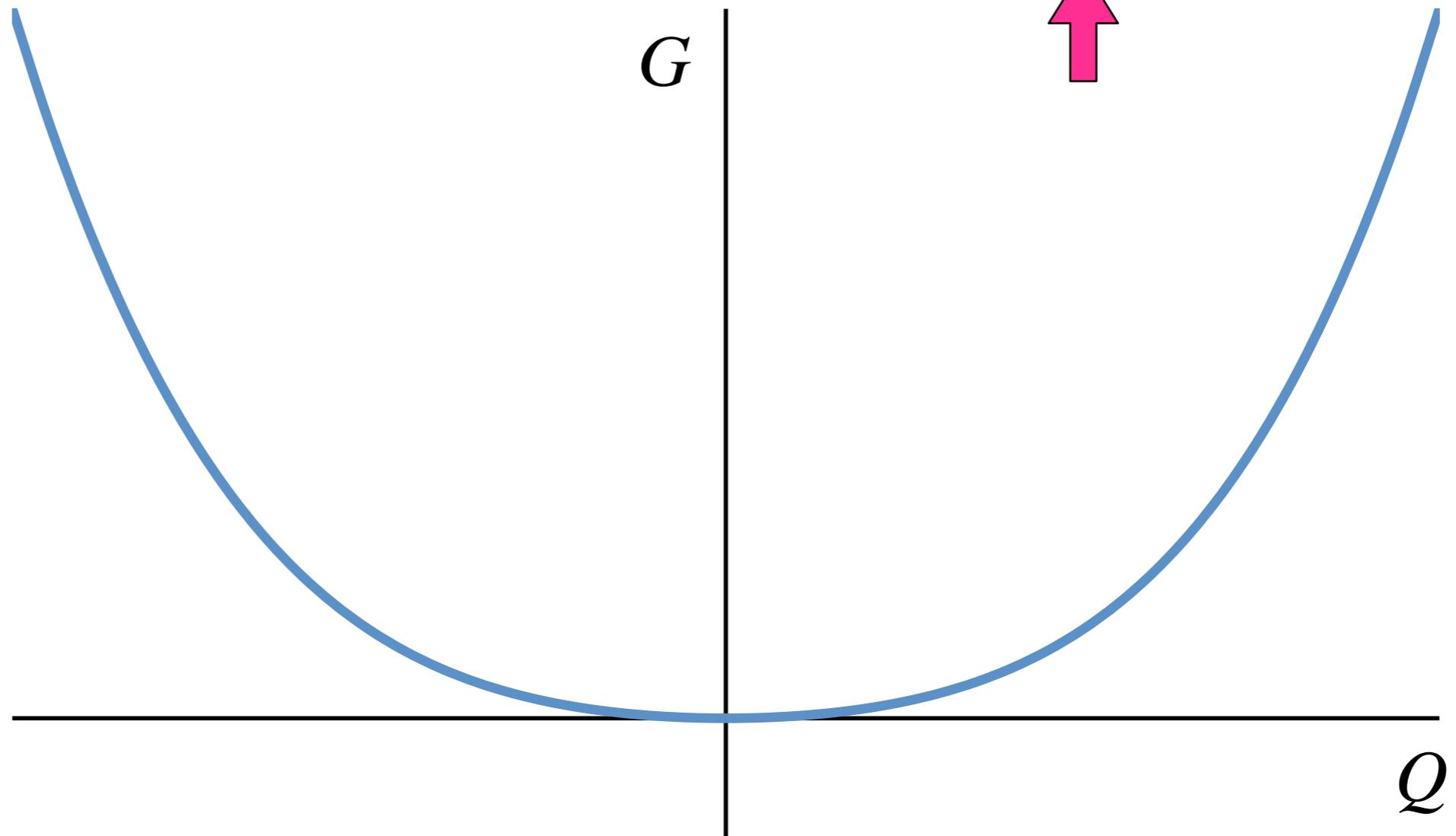
# Low temperature

$$G = G_0 - \frac{1}{2}aQ^2 + \frac{1}{4}bQ^4$$



# High temperature

$$G = G_0 + \frac{1}{2}aQ^2 + \frac{1}{4}bQ^4$$

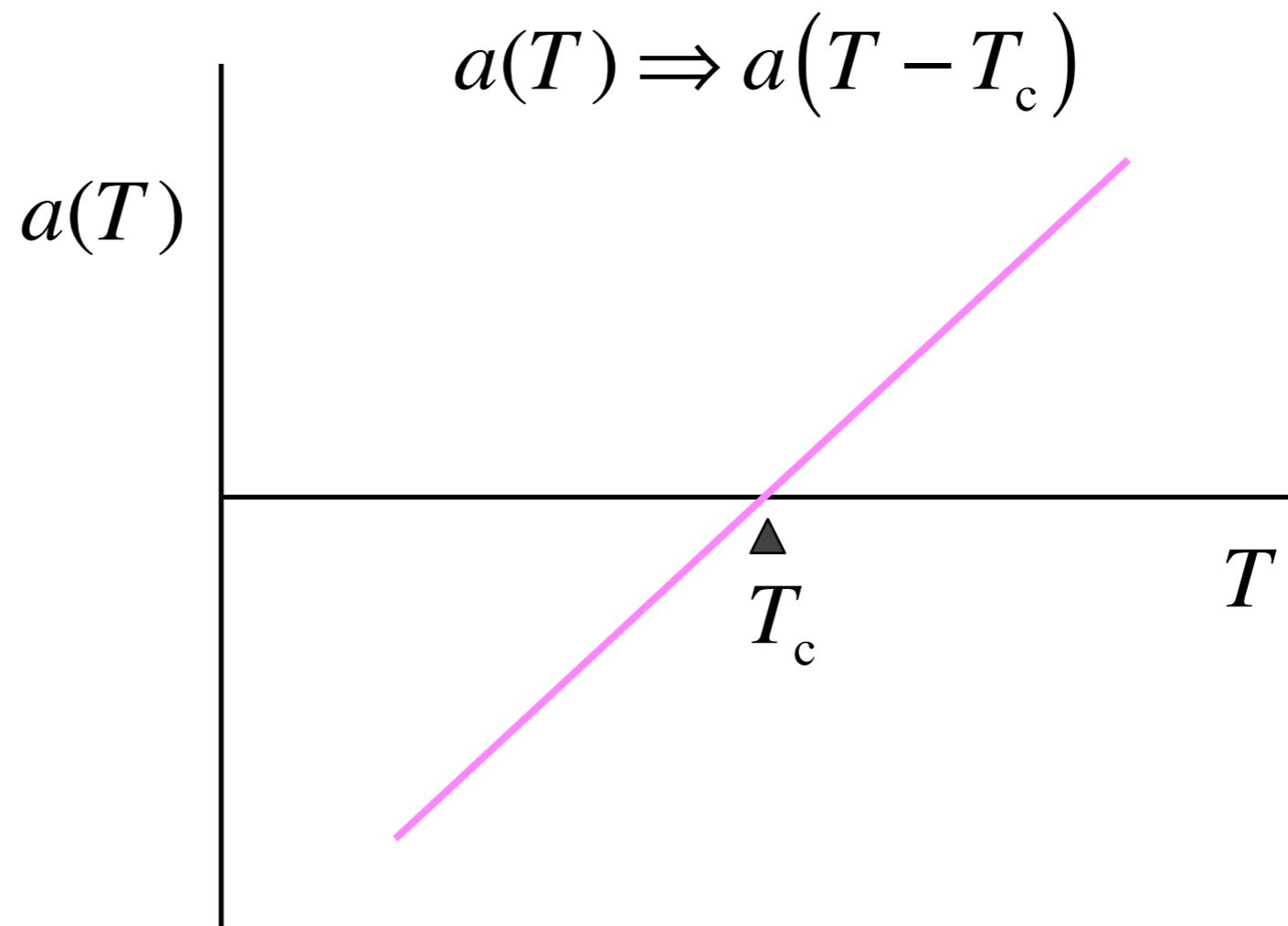


# Guess the form of $a(T)$

$$G = G_0 + \frac{1}{2}aQ^2 + \frac{1}{4}bQ^4$$

$$a > 0 \text{ for } T > T_c$$

$$a < 0 \text{ for } T < T_c$$



# Free energy function

- ▶ Final function has the form:

$$G = G_0 + \frac{1}{2}a(T - T_c)Q^2 + \frac{1}{4}bQ^4$$

- ▶ Assume the values of the coefficients are independent of temperature

# Extreme points

- ▶ Obtain the extreme point of  $G(Q)$ :

$$G = G_0 + \frac{1}{2}a(T - T_c)Q^2 + \frac{1}{4}bQ^4$$

$$\frac{\partial G}{\partial Q} = a(T - T_c)Q + bQ^3 = 0$$

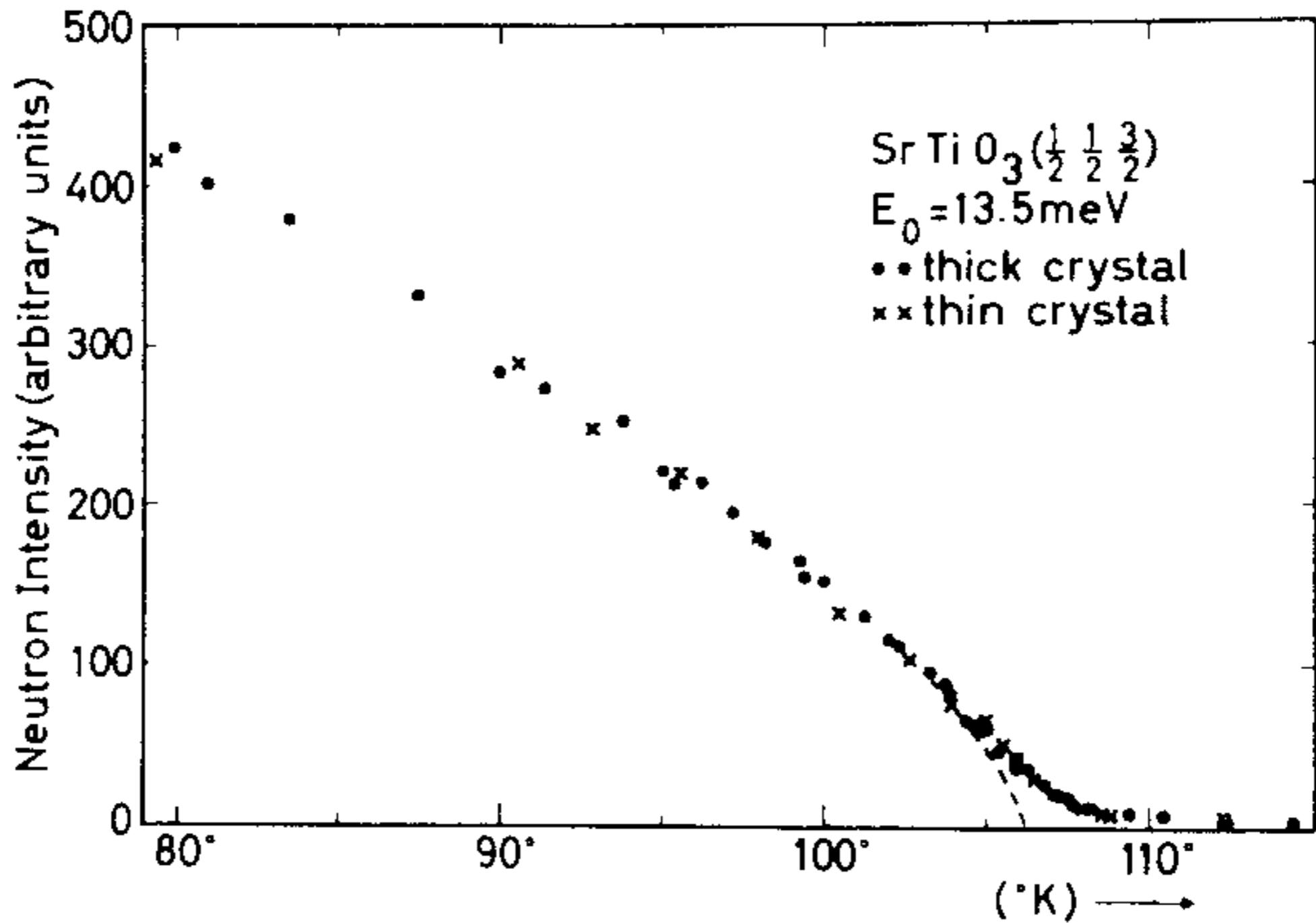
$$\Rightarrow bQ^3 = a(T_c - T)Q$$

- ▶ Two solutions:

1.  $Q = 0$

2.  $Q^2 = a(T_c - T) / b$

# Example of $\text{SrTiO}_3$



# Thermodynamic interpretation

- ▶ Separate the free energy into enthalpy and entropy terms

$$G = G_0 + \frac{1}{2}a(T - T_c)Q^2 + \frac{1}{4}bQ^4 = G_0 + \Delta H - T\Delta S$$

$$\Delta H = -\frac{1}{2}aT_cQ^2 + \frac{1}{4}bQ^4$$

$$-T\Delta S = \frac{1}{2}aTQ^2 \quad \Rightarrow \quad \Delta S = -\frac{1}{2}aQ^2$$

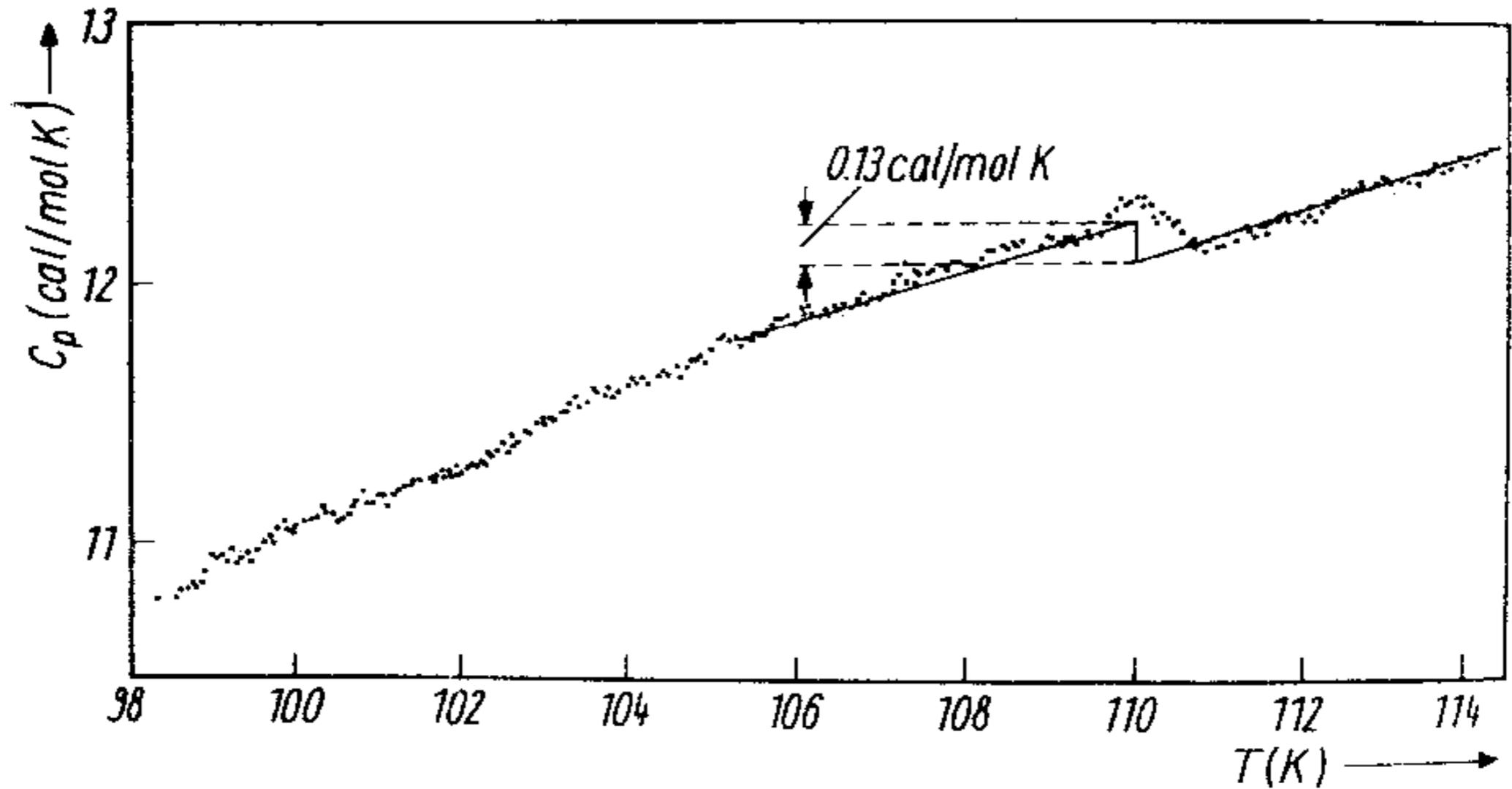
- ▶ Recall, expansion is around the high-temperature phase

# Heat capacity

$$\Delta S = -\frac{1}{2}aQ^2 = -\frac{1}{2}a^2(T_c - T)/b$$

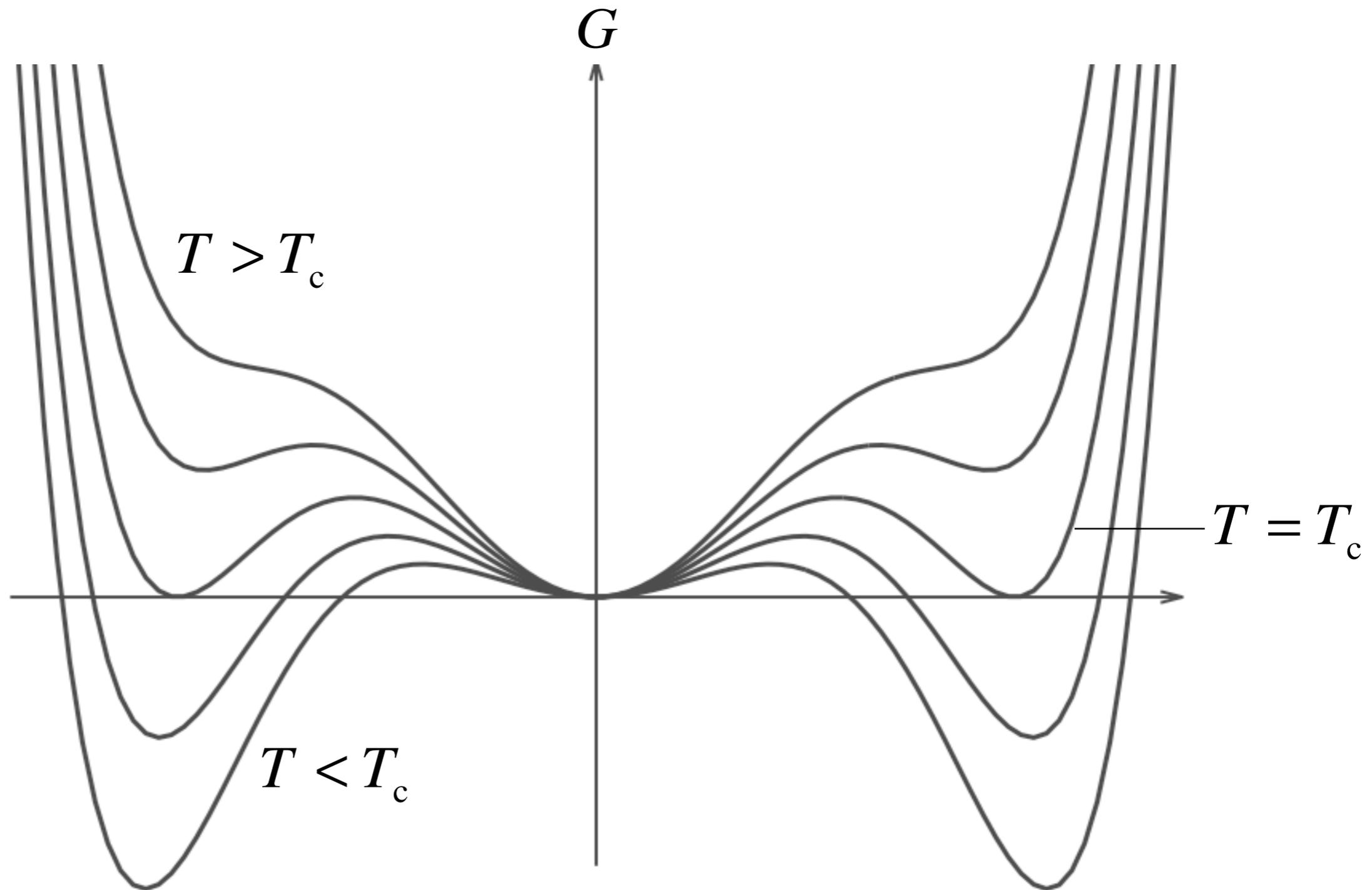
$$\Delta C = T \frac{\partial \Delta S}{\partial T} = \begin{cases} 0 & \text{for } T > T_c \\ \frac{a^2}{2b}T & \text{for } T < T_c \end{cases}$$

# Example of SrTiO<sub>3</sub>



- ▶ Small step in heat capacity at the phase transition – hard to measure

# Free energy for a first-order phase transition



# Free energy function

- ▶ Needs to be extended to sixth order with negative quartic term

$$G = G_0 + \frac{1}{2}a(T - T_0)Q^2 - \frac{1}{4}bQ^4 + \frac{1}{6}cQ^6$$

- ▶  $T_0$  is not the transition temperature

# Extreme points

$$G = G_0 + \frac{1}{2}a(T - T_0)Q^2 - \frac{1}{4}bQ^4 + \frac{1}{6}cQ^6$$

$$\frac{\partial G}{\partial Q} = a(T - T_0)Q - bQ^3 + cQ^5 = 0$$

$$\Rightarrow Q = 0 \quad \text{or} \quad Q^2 = \frac{b \pm \sqrt{b^2 - 4ac(T - T_0)}}{2c}$$

# Computing $T_c$

- ▶ At  $T = T_c$  we have

$$G - G_0 = \frac{1}{2}a(T_c - T_0)Q^2 - \frac{1}{4}bQ^4 + \frac{1}{6}cQ^6 = 0$$

$$\Rightarrow a(T_c - T_0) - \frac{1}{2}bQ^2 + \frac{1}{3}cQ^4 = 0$$

$$\frac{\partial G}{\partial Q} = a(T_c - T_0)Q - bQ^3 + cQ^5 = 0$$

$$\Rightarrow a(T_c - T_0) - bQ^2 + cQ^4 = 0$$

- ▶ The jump in the order parameter is

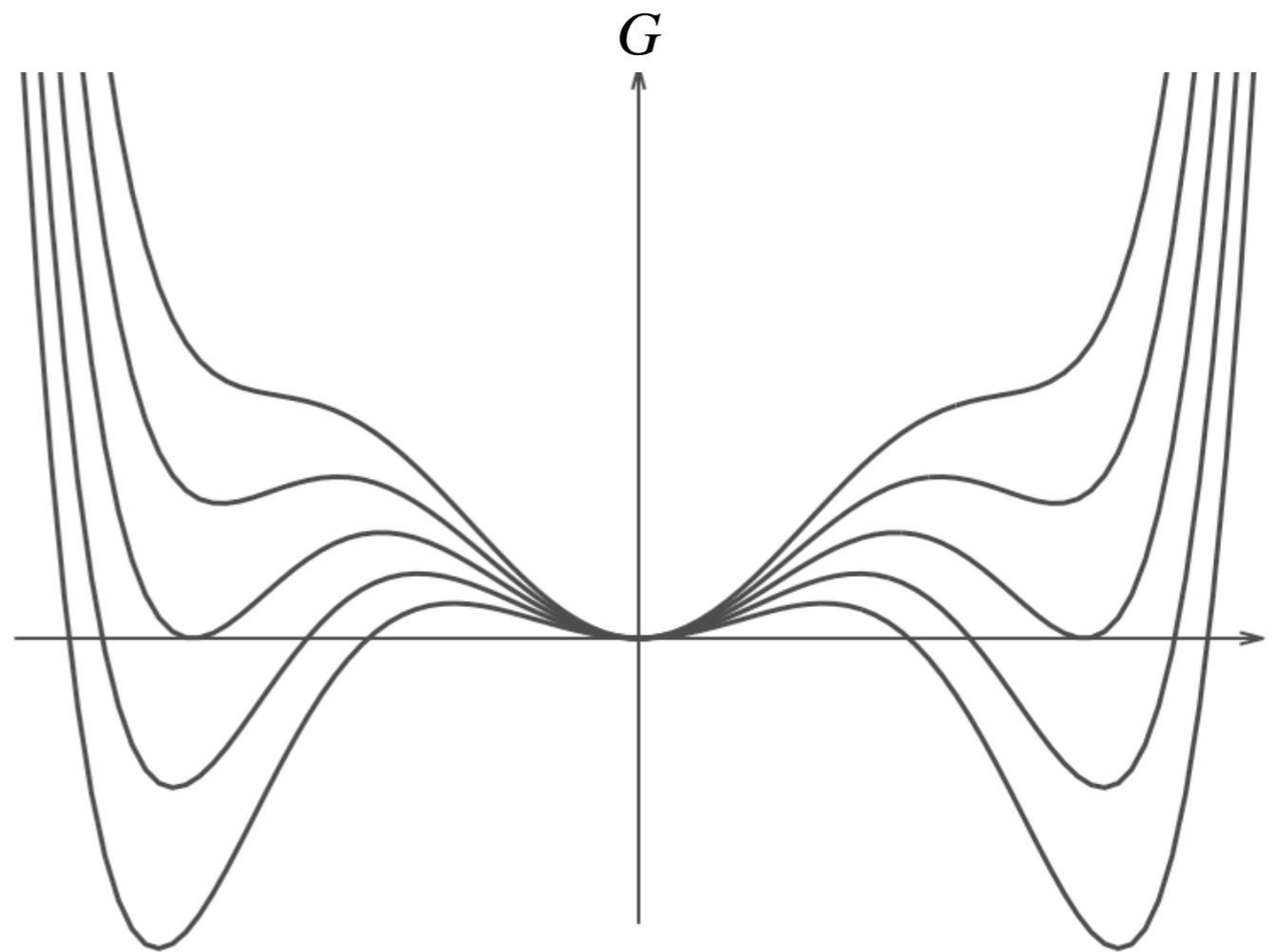
$$\Delta Q^2 = \frac{3b}{4c}$$

# Computing $T_c$

$$a(T_c - T_0) - bQ^2 + cQ^4 = 0$$

$$\Rightarrow a(T_c - T_0) - b\frac{3b}{4c} + c\left(\frac{3b}{4c}\right)^2 = 0$$

$$T_c = T_0 + \frac{3b^2}{16ac}$$



# Latent heat at $T_c$

$$\begin{aligned}\Delta H &= -\frac{1}{2}aT_0Q^2 - \frac{1}{4}bQ^4 + \frac{1}{6}cQ^6 \\ &= -\frac{3abT_c}{8c} = -\frac{aT_c}{2}\Delta Q^2\end{aligned}$$

The existence of a latent heat at a first-order phase transition is often said to be the defining feature of a first-order phase transition

# Ehrenfest categorisation

Idea is to categorise a phase transition by the first differential that is not zero

▶ First order:  $\left. \frac{\partial G}{\partial T} \right|_{T=T_c} \neq 0 \Rightarrow \Delta S, \Delta H \neq 0$

▶ Second order:  $\left. \frac{\partial^2 G}{\partial T^2} \right|_{T=T_c} \neq 0 \Rightarrow \Delta C \neq 0$

▶ Third order:  $\left. \frac{\partial^3 G}{\partial T^3} \right|_{T=T_c} \neq 0$

# Tricritical phase transitions

$$G = G_0 + \frac{1}{2}a(T - T_0)Q^2 + \frac{1}{4}bQ^4 + \frac{1}{6}cQ^6$$

$b > 0$ : Second-order

$b < 0$ : First-order

$b \sim 0$ : Tricritical

$$G = G_0 + \frac{1}{2}a(T - T_0)Q^2 + \frac{1}{6}cQ^6$$

# Order parameter

## ► Order parameter

$$G = G_0 + \frac{1}{2}a(T - T_c)Q^2 + \frac{1}{6}cQ^6$$

$$\frac{\partial G}{\partial Q} = a(T - T_c)Q + cQ^5 = 0$$

$$\Rightarrow Q = \left( \frac{a}{c}(T_c - T) \right)^{1/4}$$

## ► Entropy

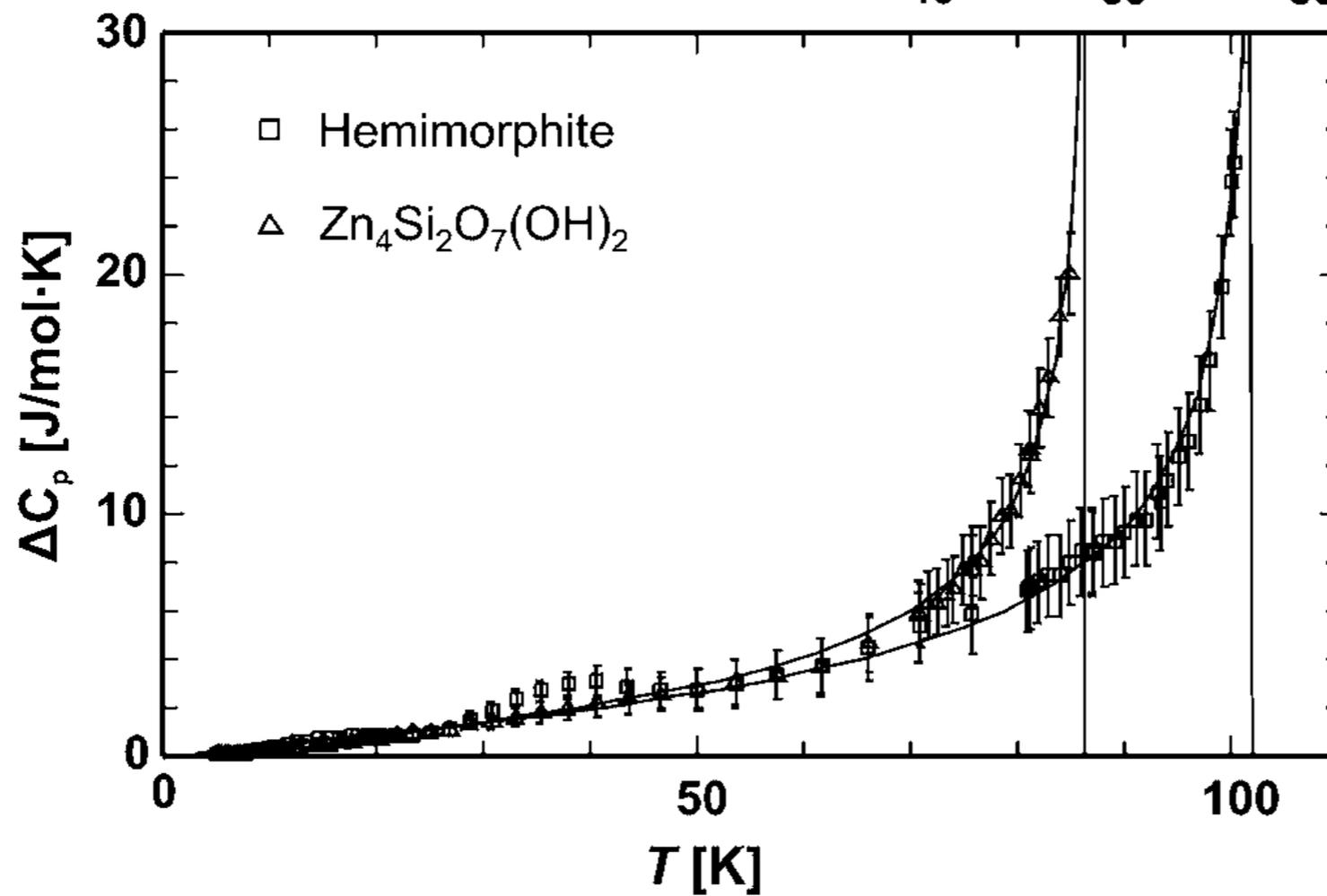
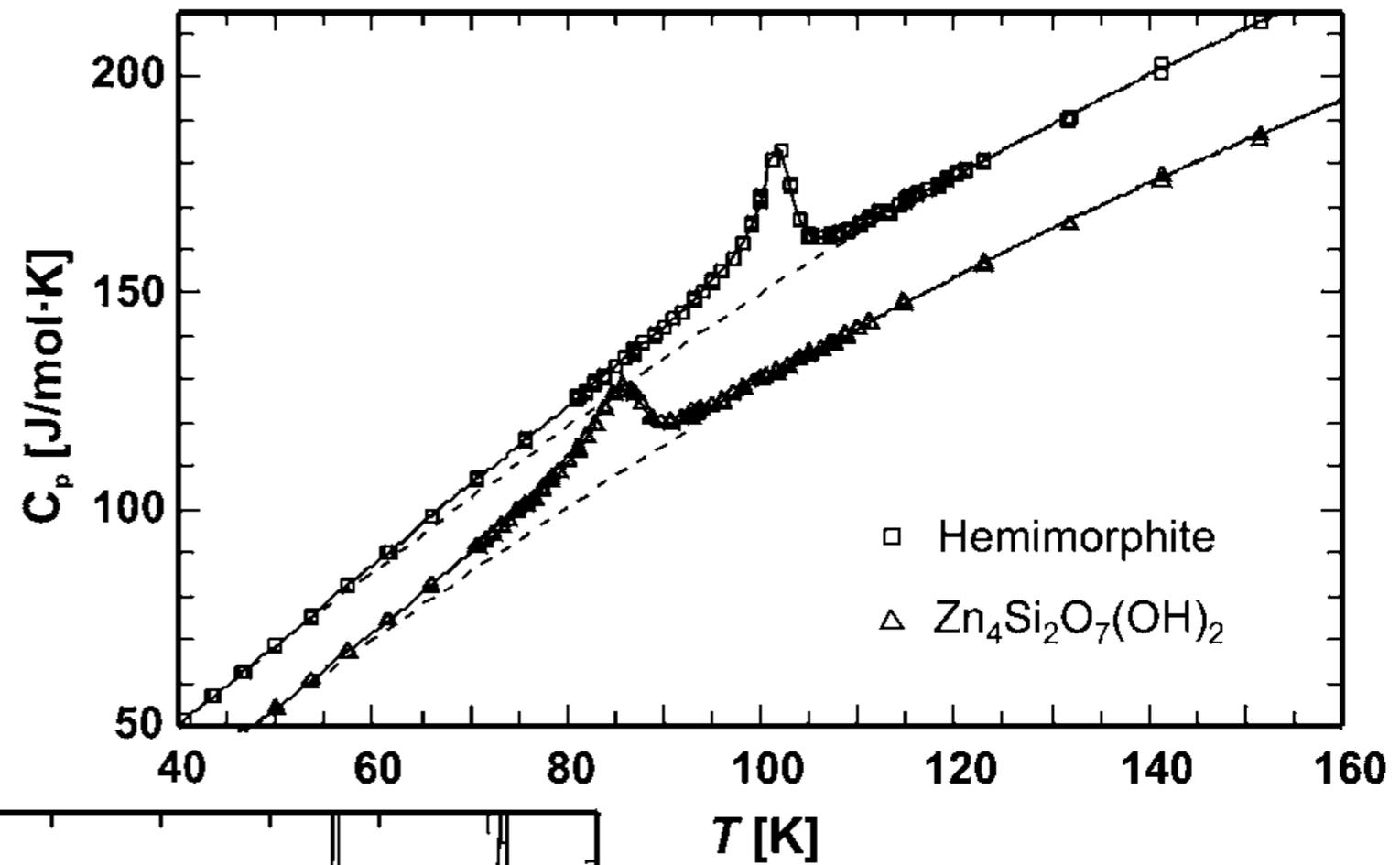
$$\Delta S = -\frac{1}{2}aQ^2 = -\frac{1}{2} \left( \frac{a^3}{c}(T_c - T) \right)^{1/2}$$

# Heat capacity

$$\Delta S = -\frac{1}{2}aQ^2 = -\frac{1}{2}\left(\frac{a^3}{c}(T_c - T)\right)^{1/2}$$

$$\Delta C = T \frac{\partial \Delta S}{\partial T} = \frac{1}{4}\sqrt{\frac{a^3}{c}}T(T_c - T)^{-1/2}$$

# Example





# Summary

- ▶ Described a free energy function that describes first and second-order phase transitions
- ▶ Seen how minimisation of the free energy leads to temperature-dependence of the order parameter
- ▶ Computed properties such as heat capacity associated with the phase transition